

Problem I

a)

$$V = \frac{\zeta Q q}{\zeta \pi \epsilon_0 s} + \frac{Q q}{\zeta \pi \epsilon_0 (d-s)}$$

$$V' = \frac{Q q}{\zeta \pi \epsilon_0} \left(-\frac{\zeta}{s^2} + \frac{1}{(d-s)^2} \right)$$

Equilibrium $\Leftrightarrow V' = 0$

$$-\frac{\zeta}{s^2} + \frac{1}{(d-s)^2} = 0$$

$$-\zeta(d-s)^2 + s^2 = 0$$

$$-\zeta d^2 + 8ds - \zeta s^2 + s^2 = 0$$

$$3s^2 - 8ds + \zeta d^2 = 0$$

$$s = \frac{8d \pm \sqrt{64d^2 - 48d^2}}{6} =$$

~~2d~~
NOT
BETWEEN
CHARGES

$$\frac{2}{3}d$$

$$\boxed{s = \frac{2}{3}d}$$

b)

$$V'' = \frac{Qq}{s\pi\epsilon_a} \left(\frac{8}{s^3} + \frac{2}{(d-s)^3} \right)$$

$$V''\left(\frac{2}{3}d\right) = \frac{Qq}{s\pi\epsilon_a} \left(\frac{8}{\frac{8}{27}d^3} + \frac{2}{\frac{1}{27}d^3} \right)$$

$$V''\left(\frac{2}{3}d\right) > 0 \text{ for stability}$$

As the term within brackets is always positive and so do Q and ϵ_a , the only possibility is

$$\boxed{q > 0}$$

so the moving charge needs to be positive for stability.

Problem II

a)

$$\underline{r}(t) = x(t) \hat{i} + y(t) \hat{j}$$

$$I(\underline{r}) = \int_0^{100} L(\underline{r}) dt$$

$$L = T - V; \quad T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2); \quad V = mgy$$

$$I(\underline{r}) = \boxed{\int_0^{100} \left(\frac{1}{2} m (\dot{x}^2 + \dot{y}^2) - mgy \right) dt}$$

b)

$$x(t) = a + bt + ct^2 \quad y(t) = d + et + ft^2$$

$$x(0) = 0 \Rightarrow \boxed{a = 0}$$

$$y(0) = 0 \Rightarrow \boxed{d = 0}$$

$$x(100) = 100 \Rightarrow 100b + 10000c = 100$$

$$b = 1 - 100c$$

$$y(100) = 10000 \Rightarrow 100e + 10000f = 10000$$

$$e = 100 - 100f$$

$$x(t) = (1 - 100c)t + ct^2$$

$$y(t) = (100 - 100f)t + ft^2$$

$$\dot{x}(t) = (1 - 100c) + 2ct$$

$$\dot{y}(t) = (100 - 100f) + 2ft$$

$$I(r) = \int_0^{100} m \left(\frac{1}{2} \left(((1 - 100c) + 2ct)^2 + ((100 - 100f) + 2ft)^2 \right) - 10((100 - 100f)t + ft^2) \right) dt$$

$$= m \int_0^{100} \left(\frac{1}{2} \left((1 - 100c)^2 + 2(1 - 100c)2ct + 4c^2t^2 + (100 - 100f)^2 + 2(100 - 100f)2ft + 4f^2t^2 \right) - 10((100 - 100f)t + ft^2) \right) dt$$

$$= m \left((1 - 100c)^2 50 + (1 - 100c) \cdot c \cdot 10000 + \frac{2}{3} c^2 \cdot 1000000 + (100 - 100f)^2 50 + (100 - 100f) \cdot f \cdot 10000 + \frac{2}{3} f^2 \cdot 1000000 - 10((100 - 100f)5000 + \frac{1}{3} f \cdot 1000000) \right)$$

$$= \Phi(c, f)$$

To find an extremal:

$$\frac{\partial \phi}{\partial c} = \frac{\partial \phi}{\partial f} = 0$$

$$\begin{aligned}\frac{\partial \phi}{\partial c} &= 100(1-100c) \cdot (-100) - 100 \cdot c \cdot 10000 + (1-100c) \cdot 10000 \\ &\quad + \frac{5}{3} c \cdot 1000000 = 0\end{aligned}$$

$$-10000 + 10000c - 1000000c + 10000 - 1000000c$$

$$+ \frac{5000000}{3} c = 0$$

$$\boxed{c = 0}$$

$$\begin{aligned}\frac{\partial \phi}{\partial f} &= 100(100-100f)(-100) - 100 \cdot f \cdot 10000 + (100-100f) \cdot 10000 \\ &\quad + \frac{5}{3} f \cdot 1000000 + 1000 \cdot 5000 - \frac{10000000}{3}\end{aligned}$$

$$-1000000 + 1000000f - 1000000f + 1000000 - 1000000f$$

$$+ \frac{5000000}{3} f + 5000000 - \frac{10000000}{3}$$

$$= \cancel{-} \frac{1000000}{3} f + \frac{5000000}{3} = 0$$

$$\boxed{f = -5}$$

$$b = 1 - 100c = 1$$

$$e = 100 - 100f = 600$$

So the solution is

$$x(t) = t$$

$$y(t) = 600t - 5t^2$$

- c) The actual path of the projectile should be a parabola, which is a consequence of the kinematics of the problem.

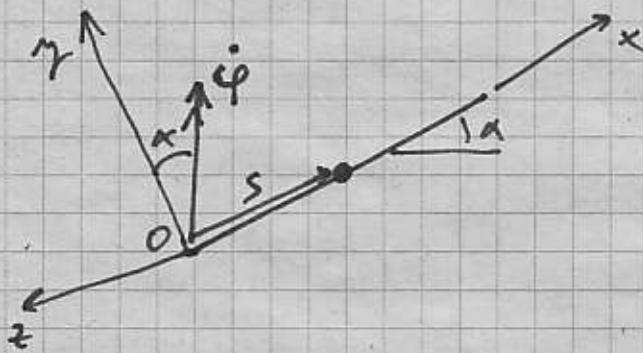
A second order expansion for Ritz method applied to Hamilton's principle provides consequently the exact answer.

This is further demonstrated by the obtained coefficients. Indeed, $c=0$ which corresponds to a linear dependence of the path on time in the x-direction.

Also the coefficient $f = -5$ corresponds to the term $-\frac{1}{2}gt^2$ of the path in the y-direction.

Problem III

a)



$$\underline{r}_{\text{rel}} = s \hat{i} \quad \underline{v}_{\text{rel}} = \dot{s} \hat{i}$$

$$\underline{\omega} = \dot{\phi} \sin \alpha \hat{i} + \dot{\phi} \cos \alpha \hat{j}$$

$$\underline{v} = \underline{v}_0 + \underline{\omega} \times \underline{r}_{\text{rel}} + \underline{v}_{\text{rel}}$$

$$= \underline{0} + \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \dot{\phi} \sin \alpha & \dot{\phi} \cos \alpha & 0 \\ s & 0 & 0 \end{vmatrix} + \dot{s} \hat{i}$$

$$= \dot{s} \hat{i} - s \dot{\phi} \cos \alpha \hat{k}$$

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) = \frac{1}{2} m (\dot{s}^2 + s^2 \dot{\phi}^2 \cos^2 \alpha)$$

$$V = mg s \sin \alpha$$

$$L = T - V = \frac{1}{2} m (\dot{s}^2 + s^2 \dot{\phi}^2 \cos^2 \alpha) - mg s \sin \alpha$$

b)

$$L \neq L(\phi) \Rightarrow \phi \text{ is ignorable}$$

$$C_\phi = \frac{\partial L}{\partial \dot{\phi}} = ms^2 \dot{\phi} \cos^2 \alpha$$

This represents conservation of angular momentum

c)

$$\frac{\partial L}{\partial \dot{s}} = m \ddot{s} \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) = m \ddot{s}$$

$$\frac{\partial L}{\partial s} = m s \dot{\varphi}^2 \cos^2 \alpha - mg \sin \alpha$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) - \frac{\partial L}{\partial s} = m \ddot{s} - m s \dot{\varphi}^2 \cos^2 \alpha + mg \sin \alpha = 0$$

$$\dot{\varphi} = \frac{C_\varphi}{m s^2 \cos^2 \alpha}$$

$$m \ddot{s} - m s \cos^2 \alpha \frac{C_\varphi^2}{m^2 s^4 \cos^4 \alpha} + mg \sin \alpha = 0$$

$$\boxed{\ddot{s} - \frac{C_\varphi^2}{m^2 s^3 \cos^2 \alpha} + g \sin \alpha = 0}$$

d)

$$\dot{s} = \ddot{s} = 0$$

$$\cancel{m s \dot{\varphi}^2 \cos^2 \alpha} = \cancel{mg \sin \alpha}$$

$$\boxed{s \dot{\varphi}^2 = \frac{g \sin \alpha}{\cos^2 \alpha}}$$

e)

$$\ddot{s} = \frac{C_\varphi^2}{m^2 s^3 \cos^2 \alpha} - g \sin \alpha$$

$$\dot{s} = r$$

$$\dot{r} = \frac{C_\varphi^2}{m^2 s^3 \cos^2 \alpha} - g \sin \alpha$$

}

Jacobian at steady motion:

$$\begin{aligned}
 & \left[\begin{array}{cc} 0 & 1 \\ -\frac{3 C_\varphi^2}{m^2 s^3 \cos^2 \alpha} & 0 \end{array} \right] \\
 & \quad s \dot{\varphi}^2 = \frac{g \sin \alpha}{\cos^2 \alpha} \\
 & = \left[\begin{array}{cc} 0 & 1 \\ -\frac{3 m^2 s^3 \dot{\varphi}^2 \cos^2 \alpha}{m^2 s^3 \cos^2 \alpha} & 0 \end{array} \right] \\
 & = \left[\begin{array}{cc} 0 & 1 \\ -\frac{3 g \sin \alpha}{s} & 0 \end{array} \right]
 \end{aligned}$$

$$\begin{vmatrix} -\gamma & 1 \\ -\frac{3g \sin \alpha}{s} & -\gamma \end{vmatrix} = \gamma^2 + \frac{3g \sin \alpha}{s} = 0$$

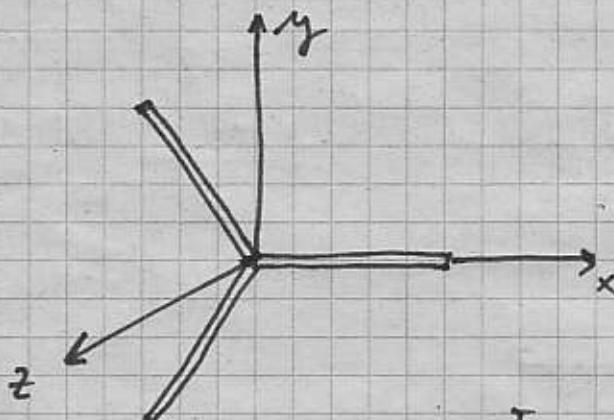
$$\gamma = \pm i \sqrt{\frac{3g \sin \alpha}{s}} \quad \text{as } s > 0 \text{ always}$$

so $\gamma \in \mathbb{C} - \mathbb{R}$

so the steady motion is stable

Problem IV

a)



$$I_3 = I_{zz} = 3 I_b$$

In xy -plane there are three symmetry axes, consequently all directions are principal in that plane, so:

$$I_{xx} = I_{yy} = I_1 = I_2 \quad (1)$$

Moreover the object is contained in the xy -plane, consequently

~~Equation~~

$$I_{zz} = \int (x^2 + y^2) dm = \int x^2 dm + \int y^2 dm \\ = I_{yy} + I_{xx} \quad (2)$$

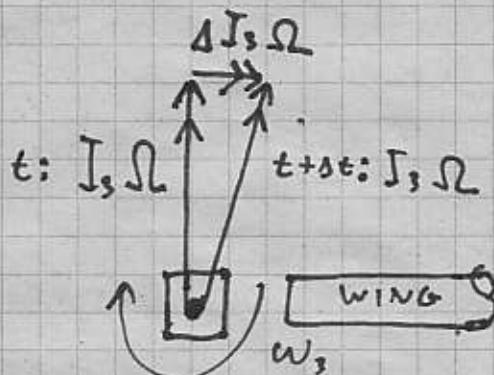
(1) and (2) together provide

$$I_{xx} = I_{yy} = \frac{I_{zz}}{2}; \quad I_1 = I_2 = \frac{I_3}{2}$$

$$I_1 = \frac{3}{2} I_b; \quad I_2 = \frac{3}{2} I_b; \quad I_3 = 3 I_b$$

b) Starboard rotor:

KD:



FBD:



So the roll causes a positive torsional moment

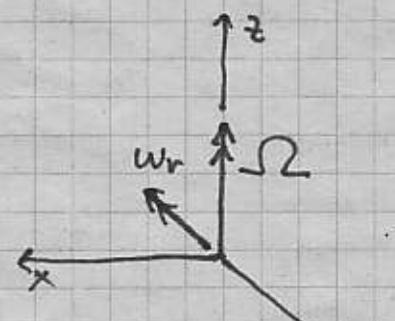


c)

$$\omega_1 = 0 \quad \omega_2 = -\omega_r \quad \omega_3 = \Omega$$

$$\underline{\omega_{12} = -\omega_r \dot{j}}$$

$$\frac{d\omega_{12}}{dt} = \dot{\omega}_1 \dot{i} + \dot{\omega}_2 \dot{j} + \begin{vmatrix} i & j & k \\ \omega_1 & \omega_2 & \omega_3 \\ \dot{\omega}_1 & \dot{\omega}_2 & 0 \end{vmatrix} = 0$$



$$\dot{\omega}_1 = \omega_2 \omega_3$$

$$\dot{\omega}_2 = -\omega_1 \omega_3$$

Euler equations:

$$\left. \begin{aligned} M_1 &= I_1 \dot{\omega}_1 - (I_2 - I_3) \omega_2 \omega_3 \\ M_2 &= I_2 \dot{\omega}_2 - (I_3 - I_1) \omega_1 \omega_3 \\ M_3 &= I_3 \dot{\omega}_3 - (I_1 - I_2) \omega_1 \omega_2 \end{aligned} \right\}$$

$$M_1 = \frac{3I_b}{2} (-\omega_r) \Omega - \left(\frac{3I_b}{2} - 3J_b \right) (-\omega_r) \Omega$$

$$= \boxed{-3I_b \omega_r \Omega = M_1}$$

$$M_2 = \frac{3I_b}{2} \cdot 0 \cdot \Omega - (3J_b - \frac{3}{2}I_b) \cdot 0 \cdot \Omega = 0$$

$$M_3 = 3I_b \cdot 0 - \left(\frac{3I_b}{2} - \frac{3I_b}{2} \right) \cdot 0 \cdot (-\omega_r) = 0$$

So $M_1 = -3I_b \omega_r \Omega$, which also
is in the sense predicted in b)