

opgave no.

naam

studienummer

vak AE3 - 915

code

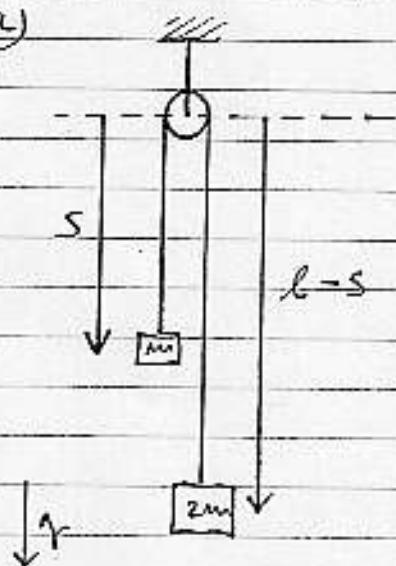
datum 27-3-08

fac.

Gebruik voor elke opgave een afzonderlijk vel papier!

1

a)

Generalised coordinate: s

$$\begin{aligned} T &= \frac{1}{2} m \dot{s}^2 + \frac{1}{2} 2m (\dot{l} - \dot{s})^2 \\ &= \frac{1}{2} m \dot{s}^2 + m \dot{l}^2 \\ &= \frac{3}{2} m \dot{s}^2 \end{aligned}$$

$$\begin{aligned} V &= -mg s - 2mg(l-s) \\ &= mg s - 2mgl \end{aligned}$$

$$L = T - V = \frac{3}{2} m \dot{s}^2 - mg s$$

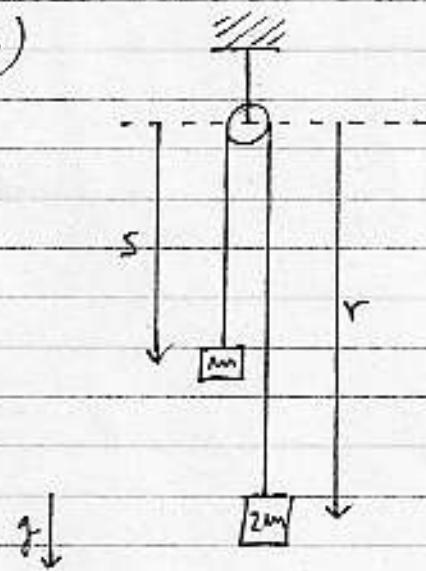
$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) - \frac{\partial L}{\partial s} = 0$$

$$\frac{\partial L}{\partial \dot{s}} = 3m\ddot{s} \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) = 3m\ddot{s} \quad \frac{\partial L}{\partial s} = -mg$$

$$3m\ddot{s} + mg = 0$$

$$\boxed{\ddot{s} + \frac{g}{3} = 0}$$

b)

Generalised coordinates: s, r

Constraint:

$$f(s, r) = s + r - l = 0$$

$$\begin{aligned} T &= \frac{1}{2} m \dot{s}^2 + \frac{1}{2} 2m \dot{r}^2 \\ &= \frac{1}{2} m \dot{s}^2 + m \dot{r}^2 \end{aligned}$$

$$V = -mg s - 2mg r$$

$$L = T - V = \frac{1}{2} m \dot{s}^2 + m \dot{r}^2 + mg s + 2mg r$$

$$\left. \begin{aligned} \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) - \frac{\partial L}{\partial s} &= \frac{\partial f}{\partial s} \\ \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) - \frac{\partial L}{\partial r} &= \frac{\partial f}{\partial r} \\ f(s, r) &= 0 \end{aligned} \right\}$$

$$\frac{\partial L}{\partial \dot{s}} = m \ddot{s} \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{s}} \right) = m \ddot{s} \quad \frac{\partial L}{\partial s} = mg \quad \frac{\partial f}{\partial s} = 1$$

$$\frac{\partial L}{\partial \dot{r}} = 2m \ddot{r} \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) = 2m \ddot{r} \quad \frac{\partial L}{\partial r} = 2mg \quad \frac{\partial f}{\partial r} = 1$$

$$\begin{aligned} m\ddot{s} - mg &= \ddot{a} \\ 2m\ddot{r} - 2mg &= \ddot{a} \\ s + r &= l \end{aligned} \quad \left. \begin{array}{l} \\ \\ \end{array} \right\}$$

Problem 2

General expression for velocity.

$$\underline{v} = \underline{v}_{xyz} + \underline{\omega} \times \underline{r}_{rel} + \underline{v}_{rel}$$

Rigid Body: $\Rightarrow \underline{v}_{rel} = 0$

Origin fixed: $\underline{v}_{xyz} = 0$ $\underline{r}_{rel} = \underline{r}$

$$\underline{v} = \underline{\omega} \times \underline{r}$$

$$T = \frac{1}{2} \int_V \underline{v} \cdot \underline{v} dm = \frac{1}{2} \int_V (\underline{\omega} \times \underline{r}) \cdot (\underline{\omega} \times \underline{r}) dm$$

$$= \frac{1}{2} \int_V \begin{vmatrix} i & j & k \\ \omega_x & \omega_y & \omega_z \\ x & y & z \end{vmatrix}^2 dm$$

$$= \frac{1}{2} \int_V [(\omega_y z - \omega_z y) \hat{i} + (\omega_z x - \omega_x z) \hat{j} + (\omega_x y - \omega_y x) \hat{k}]^2 dm$$

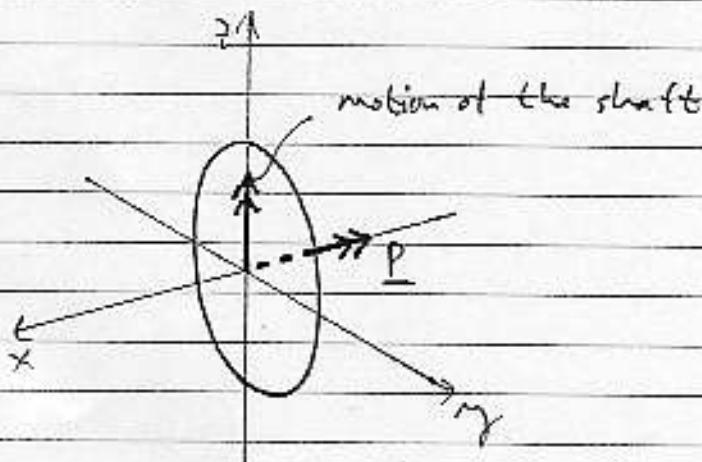
$$= \frac{1}{2} \int_V [(y^2 + z^2) \omega_x^2 + (x^2 + z^2) \omega_y^2 + (x^2 + y^2) \omega_z^2 - 2xy \omega_x \omega_y - 2xz \omega_x \omega_z - 2yz \omega_y \omega_z] dm$$

$$= \frac{1}{2} (\omega_x \ \omega_y \ \omega_z) \int_V \begin{bmatrix} y^2 + z^2 & -xy & -xz \\ -xy & x^2 + z^2 & -yz \\ -xz & -yz & x^2 + y^2 \end{bmatrix} dm \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix}$$

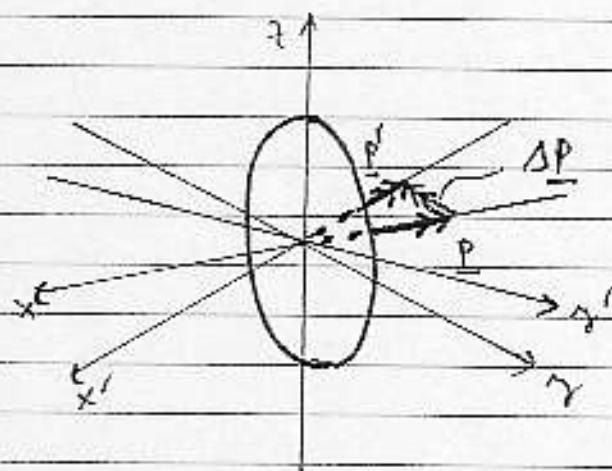
$$= \frac{1}{2} \underline{\omega}^T \underline{I}_o \underline{\omega}$$

3

Angular momentum diagram:

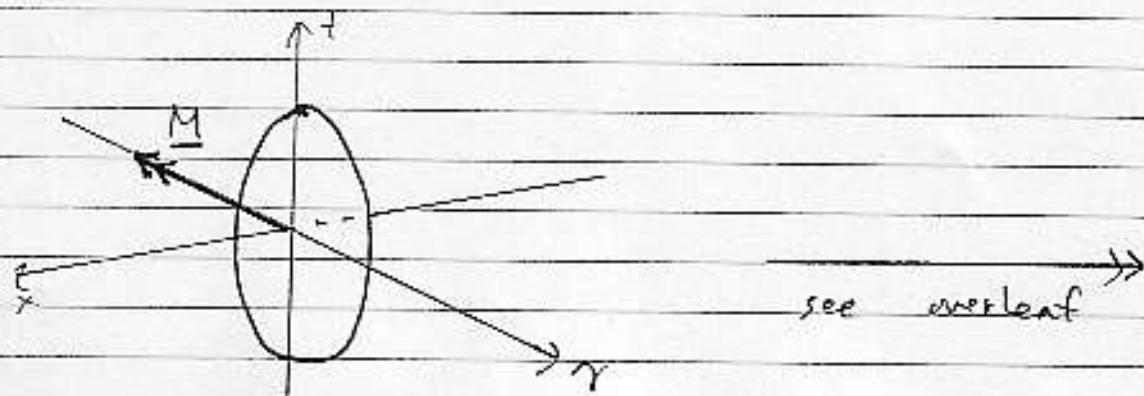


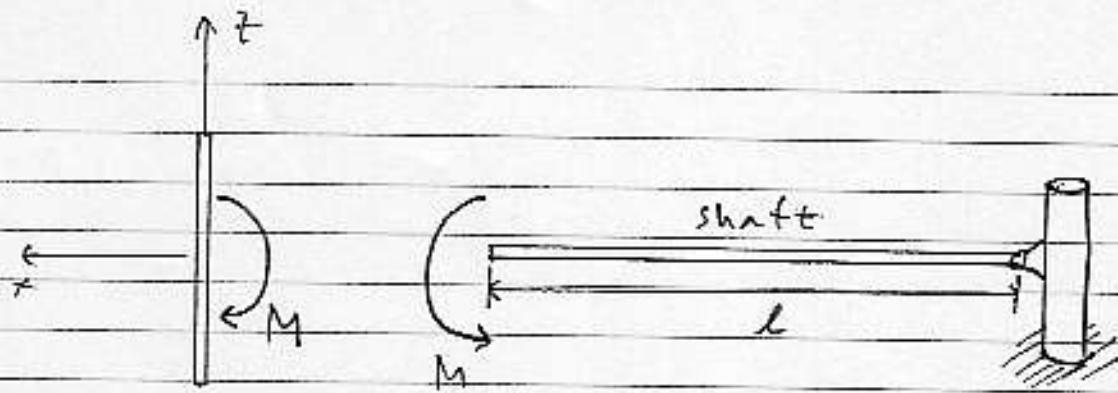
Change in angular momentum due to motion of the shaft:



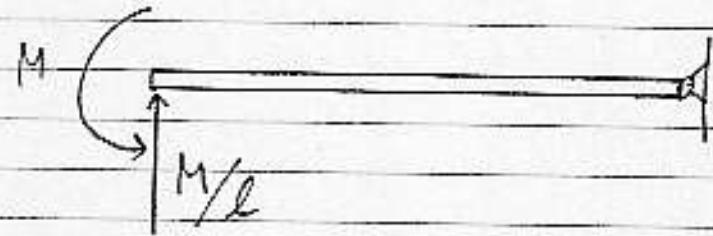
When $\Delta t \rightarrow 0$ $\Delta \underline{P}$ is parallel to the y -axis.

The corresponding FBD on the disk is:

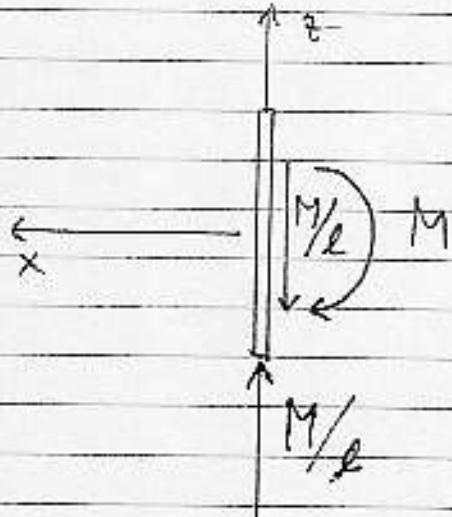


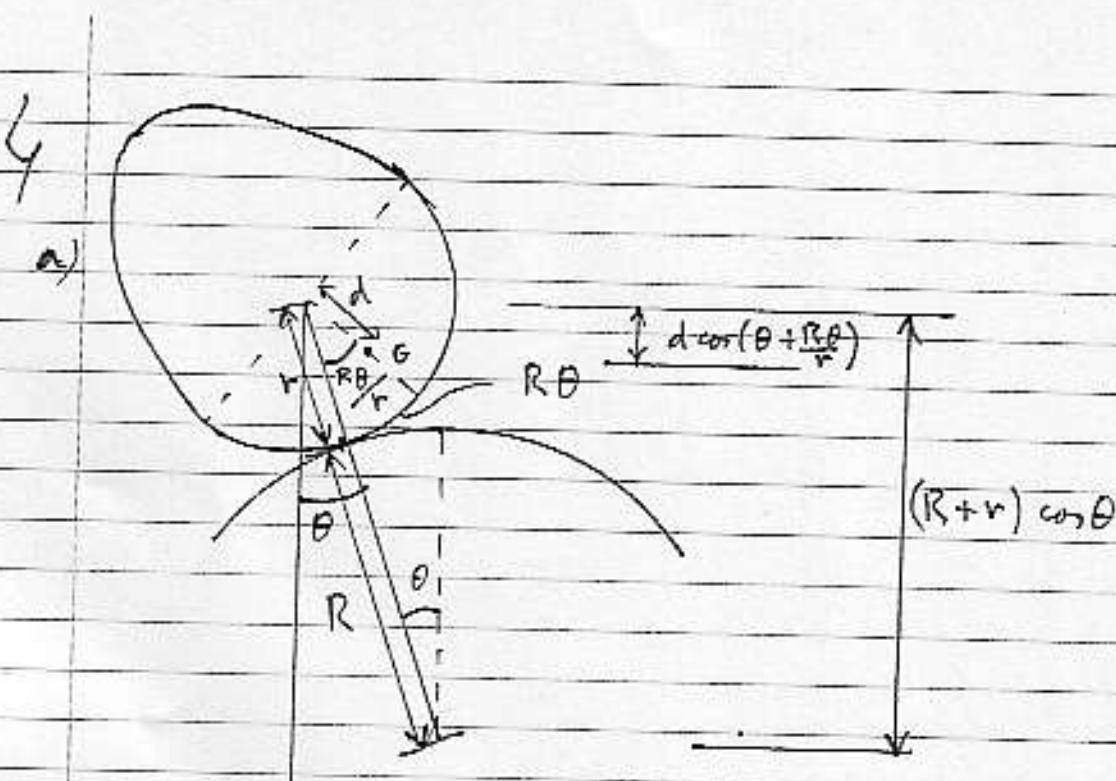


This corresponds to a CCW moment applied to the shaft. This needs to be equilibrated by an upward reaction force amounting to M/l , with l being the length of the shaft



Consequently, the static reaction at the contact point is increased by M/l





$$V = mg \left[(R+r)\cos\theta - d\cos\left(\frac{R+r}{r}\theta\right) \right]$$

b) $R = 0,012 \text{ m}$, $r = 0,021 \text{ m}$, $m = 0,05 \text{ kg}$, $g = 10 \text{ m/s}^2$

$$V = 0,5 \left(0,063 \cos\theta - d \cos 3\theta \right)$$

$$V' = 0,5 (-0,063 \sin\theta + 3d \sin 3\theta)$$

$$V'(0) = 0 \Rightarrow \theta \text{ is equilibrium point}$$

$$V'' = 0,5 (-0,063 \cos\theta + 9d \cos 3\theta)$$

$$V''(0) = 0,5 (-0,063 + 9d) > 0 \text{ for stability}$$

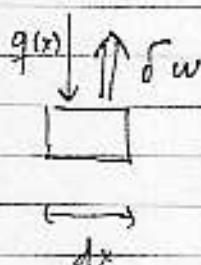
$$\boxed{d > 0,007 \text{ m}}$$

5

a)

The distributed load provides a virtual work density δW^e as

$$\delta W^e = -q(x) \delta w = -c w_x \delta w$$



It is not immediate to devise a generalized potential density that would lead to that virtual work density upon taking a variation.

An experienced analyst would propose the expression

$$\mathcal{V}^e = -\frac{1}{2} c w_x^2$$

after stating an analogy between the second spatial derivative vs. the squared first derivative, and the acceleration vs. the kinetic energy on a single-degree-of-freedom system.

When one is able to find \mathcal{V}^e the virtual work density is incorporated to the variation of the Lagrangian when applying Hamilton's principle, see answer to question C.

b)

Kinetic energy density:

$$\mathcal{T} = \frac{1}{2} \rho A u_z^2$$

Elastic potential energy density:

$$V^e = \frac{1}{2} EI w_{xx}^2$$

Lagrangian density:

$$L = \mathcal{T} - V^e = \frac{1}{2} \rho A u_z^2 - \frac{1}{2} EI w_{xx}^2$$

Remark:

If the analyst has been capable of finding a generalised potential density as shown in question a) then this should be added to the Lagrangian density as

$$L = \mathcal{T} - V^e - V^f = \frac{1}{2} \rho A u_z^2 - \frac{1}{2} EI w_{xx}^2 + \frac{1}{2} C u_x^2$$

5)

Consider a time interval $(t_1, t_2]$, the motion is such that the action is stationary:

$$\delta I(w) = \int_{t_1}^{t_2} (\delta L + \delta W) dt = 0$$

Remark: if a generalised potential has been found this is incorporated into L and then $\delta W = 0$

$$\delta I(w) = \int_{t_1}^{t_2} \int_0^L (\delta L + \delta W) dx dt = 0$$

$$= \int_{t_1}^{t_2} \int_0^L \left\{ \delta \left[\frac{1}{2} \rho A w_t^2 - \frac{1}{2} E J w_{xx}^2 \right] + [-c w_{xx} \delta w] \right\} dx dt$$

$$= \int_{t_1}^{t_2} \int_0^L (\delta L(w_t, w_{xx}) - c w_{xx} \delta w) dx dt$$

$$= \int_{t_1}^{t_2} \int_0^L \left(\frac{\partial L}{\partial w_t} \delta w_t + \frac{\partial L}{\partial w_{xx}} \delta w_{xx} - c w_{xx} \delta w \right) dx dt$$

$$= \int_0^{t_2} \left[\frac{\partial L}{\partial w_t} \delta w \right]_{t_1}^{t_2} - \int_{t_1}^{t_2} \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial w_t} \right) \delta w dt +$$

$$+ \int_{t_1}^{t_2} \left[\frac{\partial L}{\partial w_{xx}} \delta w_x \right]_0^L - \int_0^L \frac{\partial}{\partial x} \left(\frac{\partial L}{\partial w_{xx}} \right) \delta w_x dx dt + \iint_{t_1}^{t_2} -c w_{xx} \delta w dx dt$$

$$= \int_0^l \left[\frac{\partial L}{\partial w_t} \delta w \right]_{t_1}^{t_2} dx - \int_0^l \int_{t_1}^{t_2} \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial w_t} \right) \delta w dt dx$$

$$+ \int_{t_1}^{t_2} \left[\frac{\partial L}{\partial w_{xx}} \delta w_x \right]_0^l dt - \int_{t_1}^{t_2} \frac{\partial}{\partial x} \left(\frac{\partial L}{\partial w_{xx}} \right) \delta w \Big|_0^l dt$$

$$+ \int_{t_1}^{t_2} \int_0^l \frac{\partial^2}{\partial x^2} \left(\frac{\partial L}{\partial w_{xx}} \right) \delta w dx dt - \int_{t_1}^{t_2} \int_0^l c w_{xx} \delta w dx dt$$

$$= \int_{t_1}^{t_2} \int_0^l \left[\frac{\partial^2}{\partial x^2} \left(\frac{\partial L}{\partial w_{xx}} \right) - \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial w_t} \right) - c w_{xx} \right] \delta w dx dt$$

+ BOUNDARY TERMS (NOT CONSIDERED IN
THIS PROBLEM)

$$\delta I = 0 \text{ for any } \delta w \Rightarrow$$

$$\frac{\partial^2}{\partial x^2} \left(\frac{\partial L}{\partial w_{xx}} \right) - c w_{xx} - \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial w_t} \right) = 0$$

$$\boxed{\rho A w_{tt} + c w_{xx} + E J w_{xxxx} = 0}$$

EQUATION OF MOTION