

# Fluid mechanics (wb1225)

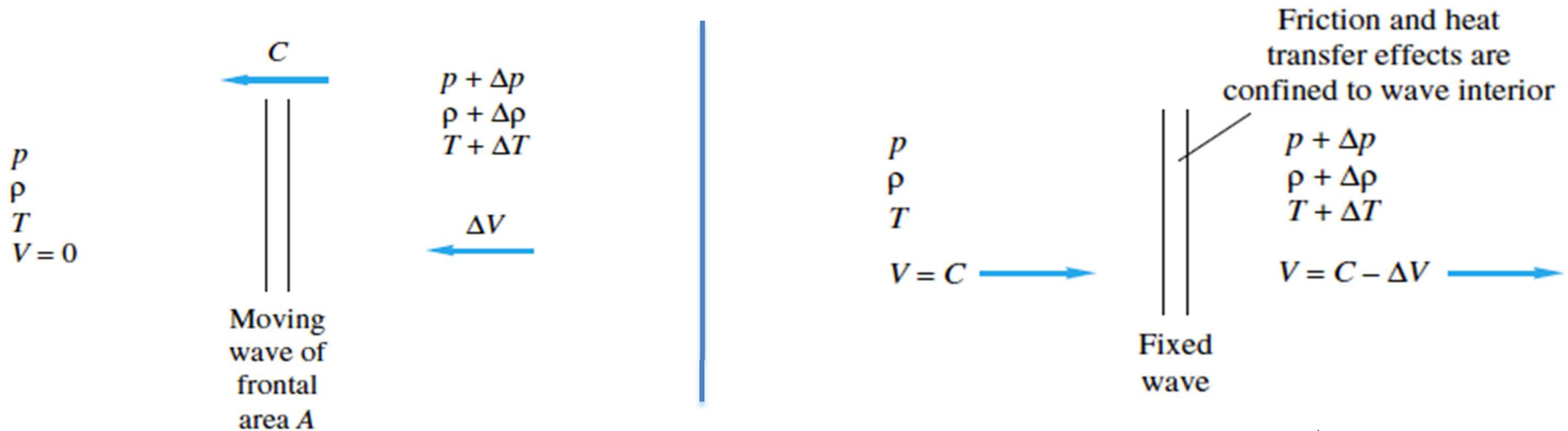
## Lecture 12: compressible flow

# Example



[1]

# Pressure wave



conservation of mass (continuity):  $\rho AC = (\rho + \Delta\rho)A(C - \Delta V) \Rightarrow \Delta V = C \frac{\Delta\rho}{\rho + \Delta\rho}$

momentum balance:

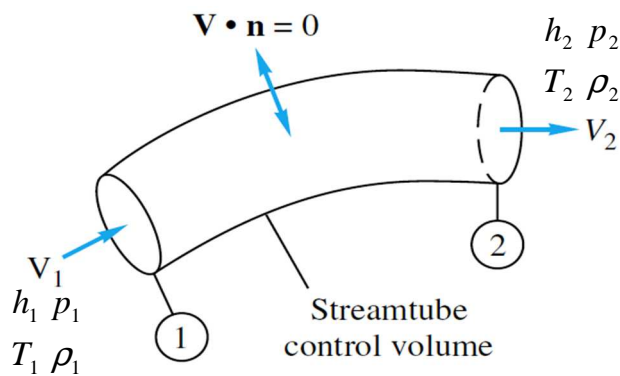
$$F = \dot{m} (V_{\text{out}} - V_{\text{in}}) \Rightarrow pA - (p + \Delta p)A = (\rho AC)(C - \Delta V - C) \Rightarrow \Delta p = \rho C \Delta V$$

combine:

$$C^2 = \frac{\Delta p}{\Delta \rho} \left( 1 + \frac{\Delta \rho}{\rho} \right) \xrightarrow{\Delta \rho / \rho \rightarrow 0} a^2 = \left. \frac{\partial p}{\partial \rho} \right|_s = kRT \begin{cases} p = \rho RT & \text{ideal gas} \\ p \rho^{-k} = \text{const} & \text{adiab. proc.} \end{cases}$$

# Adiabatic gas flow

streamtube



Conservation of energy for stationary flow under adiabatic conditions and without technical work:

- no heat transfer
- effect of viscosity occurs where  $V = 0$  (no technical work)

ideal gas

$$h = c_p T$$

$$h_1 + \frac{1}{2} V_1^2 + \cancel{gz_1} = h_2 + \frac{1}{2} V_2^2 + \cancel{gz_2}$$

changes in potential energy are negligible

$$c_p T + \frac{1}{2} V^2 = \text{constant}$$

also valid in case of losses

$$\left. \begin{aligned} c_p T + \frac{1}{2} V^2 &= c_p T_0 \\ a^2 &= \kappa R T = (\kappa - 1) c_p T \end{aligned} \right\} \Rightarrow$$

$$1 + \frac{\kappa - 1}{2} \text{Ma}^2 = \frac{T_0}{T}$$

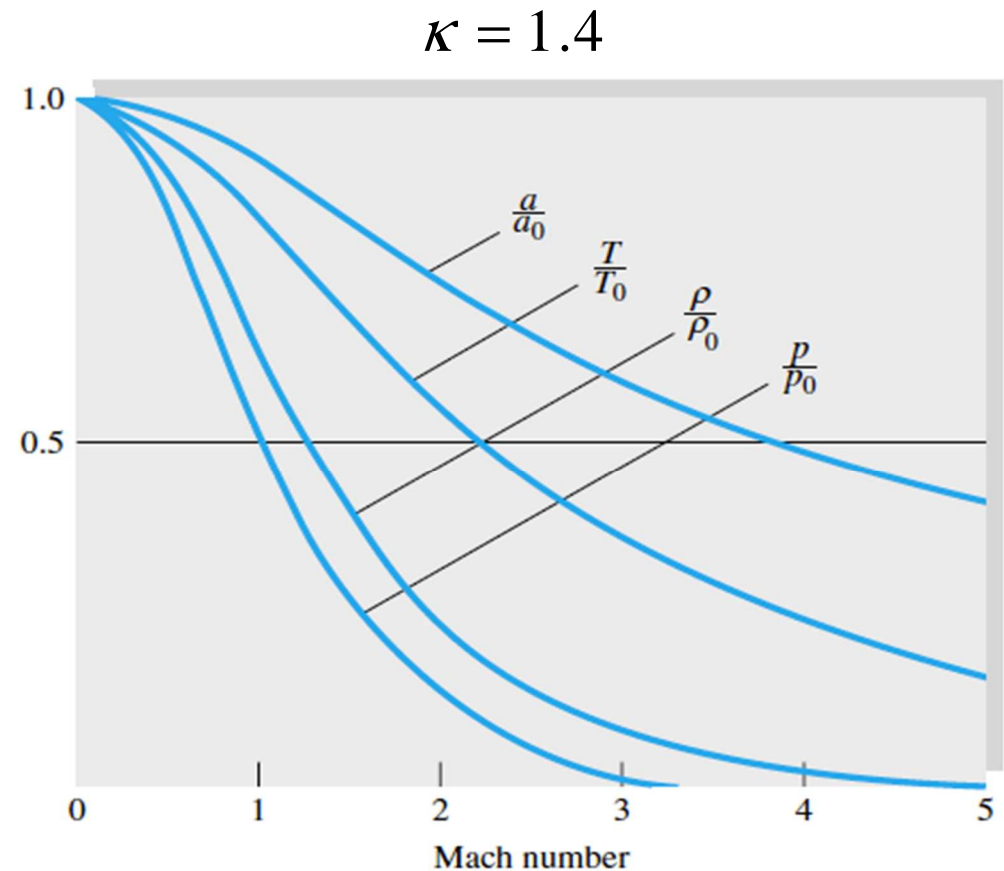
$$\text{Ma} = \frac{V}{a}$$

# Adiabatic gas flow

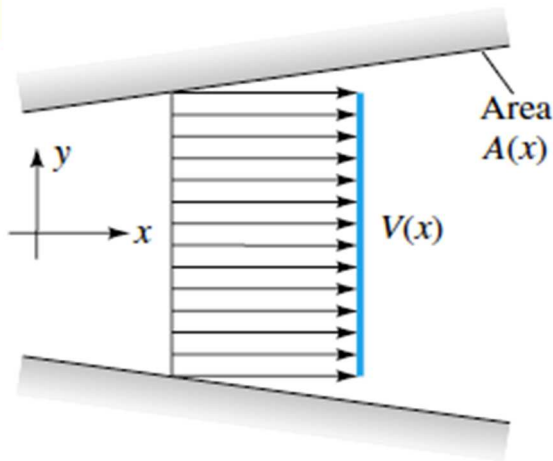
$$\frac{T_0}{T} = 1 + \frac{\kappa - 1}{2} \text{Ma}^2$$

$$\frac{p_0}{p} = \left( \frac{T_0}{T} \right)^{\frac{\kappa}{\kappa - 1}} = \left[ 1 + \frac{\kappa - 1}{2} \text{Ma}^2 \right]^{\frac{\kappa}{\kappa - 1}}$$

$$\frac{\rho_0}{\rho} = \left( \frac{T_0}{T} \right)^{\frac{1}{\kappa - 1}} = \left[ 1 + \frac{\kappa - 1}{2} \text{Ma}^2 \right]^{\frac{1}{\kappa - 1}}$$



# One-dimensional isentropic flow



continuity:

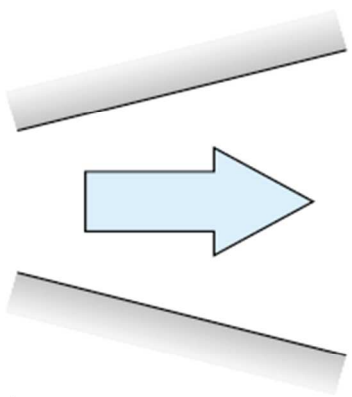
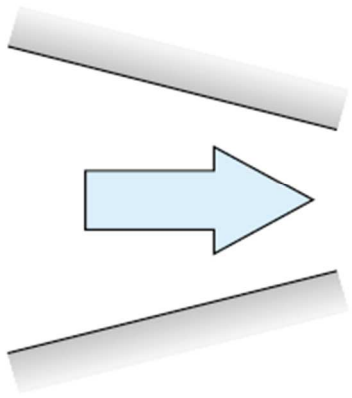
$$\dot{m} = \rho(x)V(x)A(x) = \text{const.} \quad \Rightarrow \quad \frac{d\rho}{\rho} + \frac{dV}{V} + \frac{dA}{A} = 0$$

$$\text{momentum: } \frac{dp}{\rho} + VdV = 0$$

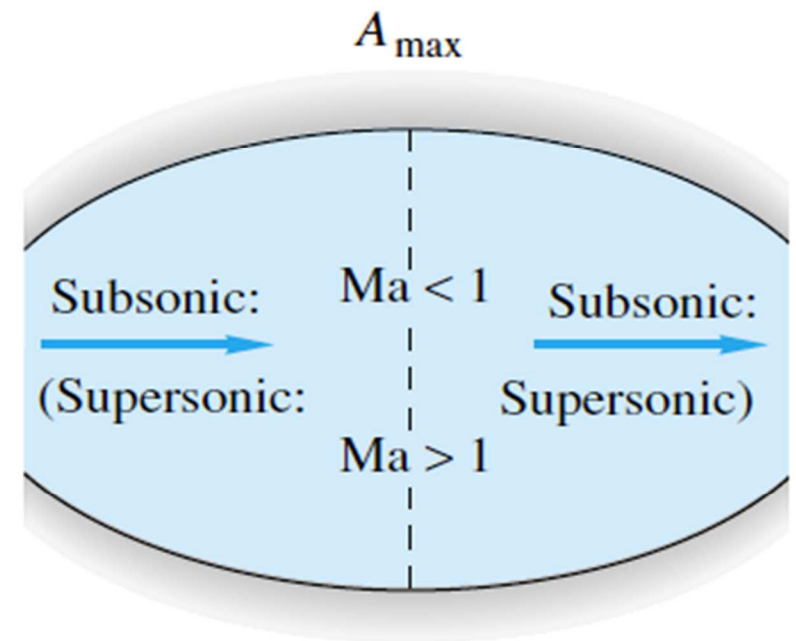
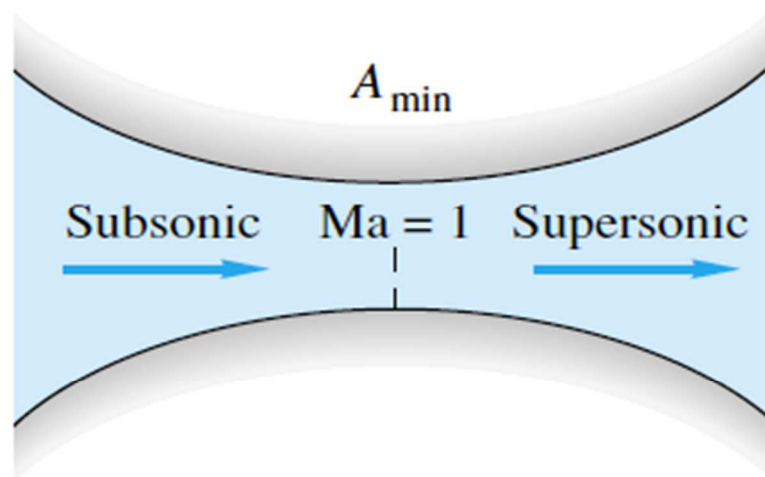
$$\text{sound: } dp = a^2 d\rho$$

$$\frac{dV}{V} = \frac{dA}{A} \frac{1}{\text{Ma}^2 - 1} = -\frac{dp}{\rho V^2}$$

# One-dimensional isentropic flow

<i>Duct geometry</i>	<i>Subsonic Ma &lt; 1</i>	<i>Supersonic Ma &gt; 1</i>
	$dA > 0$ $dV < 0$ $dp > 0$ Subsonic diffuser	$dV > 0$ $dp < 0$ Supersonic nozzle
	$dA < 0$ $dV > 0$ $dp < 0$ Subsonic nozzle	$dV < 0$ $dp > 0$ Supersonic diffuser

# Sonic flow through a nozzle





# Ideal gas flow

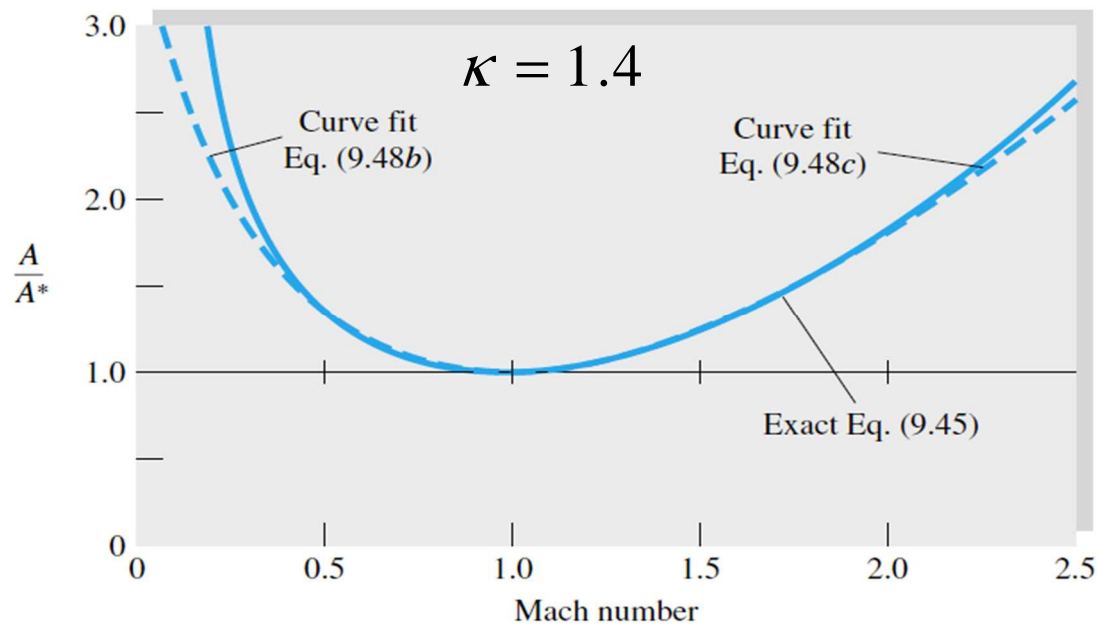
$$\rho VA = \rho^* V^* A^* \quad (\text{sonic conditions, } \text{Ma} = 1) \quad \Rightarrow \quad \frac{A}{A^*} = \frac{\rho^*}{\rho} \frac{V^*}{V}$$

$$\frac{\rho^*}{\rho} = \frac{\rho^*}{\rho_0} \frac{\rho_0}{\rho} = \frac{\left[1 + \frac{\kappa - 1}{2} \text{Ma}^2\right]^{\frac{1}{\kappa - 1}}}{\left[1 + \frac{\kappa - 1}{2} \cdot 1^2\right]^{\frac{1}{\kappa - 1}}} = \left\{ \frac{2}{\kappa + 1} \left[1 + \frac{\kappa - 1}{2} \text{Ma}^2\right] \right\}^{\frac{1}{\kappa - 1}}$$

$$\frac{V^*}{V} = \frac{(\kappa RT^*)^{1/2}}{V} = \frac{(\kappa RT)^{1/2}}{V} \left(\frac{T^*}{T}\right)^{1/2} \left(\frac{T_0}{T}\right)^{1/2} = \frac{1}{\text{Ma}} \left\{ \frac{2}{\kappa + 1} \left[1 + \frac{\kappa - 1}{2} \text{Ma}^2\right] \right\}^{\frac{1}{2}}$$

$$\frac{A}{A^*} = \frac{1}{\text{Ma}} \left[ \frac{1 + \frac{1}{2}(\kappa - 1)\text{Ma}^2}{\frac{1}{2}(\kappa + 1)} \right]^{\frac{1}{2} \frac{\kappa + 1}{\kappa - 1}}$$

# Choking



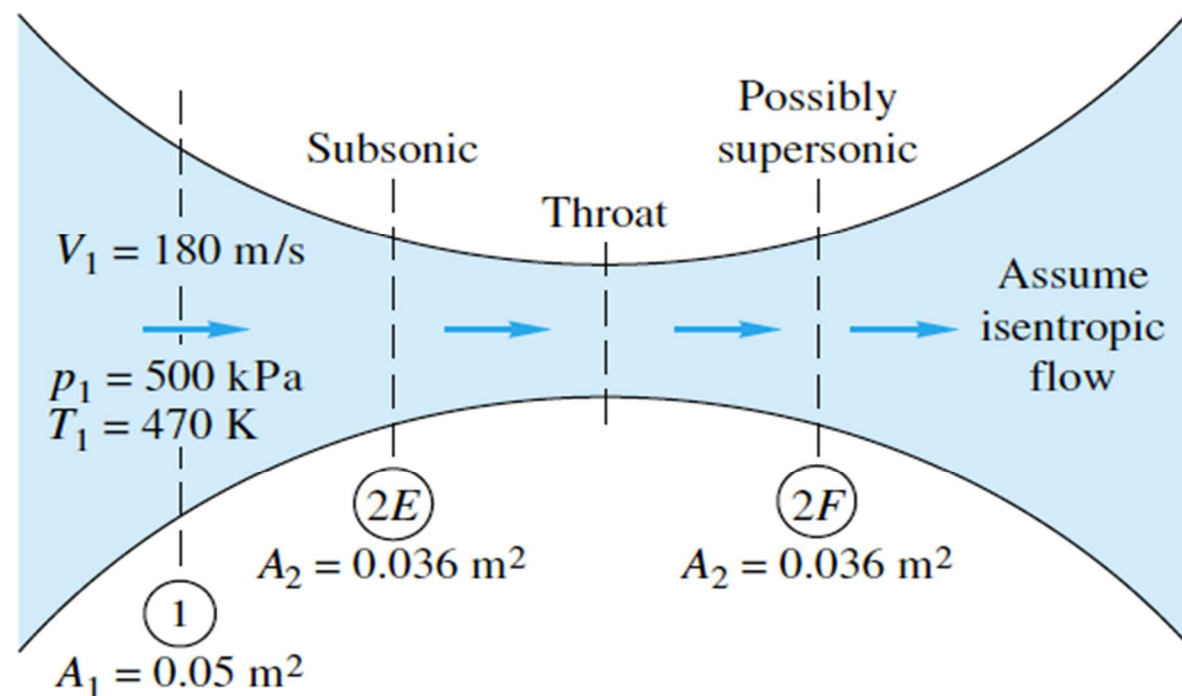
$$\frac{A}{A^*} = \frac{1}{\text{Ma}} \left[ \frac{1 + \frac{1}{2}(\kappa - 1)\text{Ma}^2}{\frac{1}{2}(\kappa + 1)} \right]^{\frac{1}{2} \frac{\kappa + 1}{\kappa - 1}}$$

$$\begin{aligned} \dot{m}_{\max} &= \rho^* A^* V^* = \rho_0 \left( \frac{2}{\kappa - 1} \right)^{\frac{1}{\kappa - 1}} A^* \left( \frac{2\kappa}{\kappa + 1} RT_0 \right)^{\frac{1}{2}} \\ &= \kappa^{\frac{1}{2}} \left( \frac{2}{\kappa + 1} \right)^{\frac{1}{2} \frac{\kappa + 1}{\kappa - 1}} A^* \rho_0 (RT_0)^{\frac{1}{2}} \end{aligned}$$

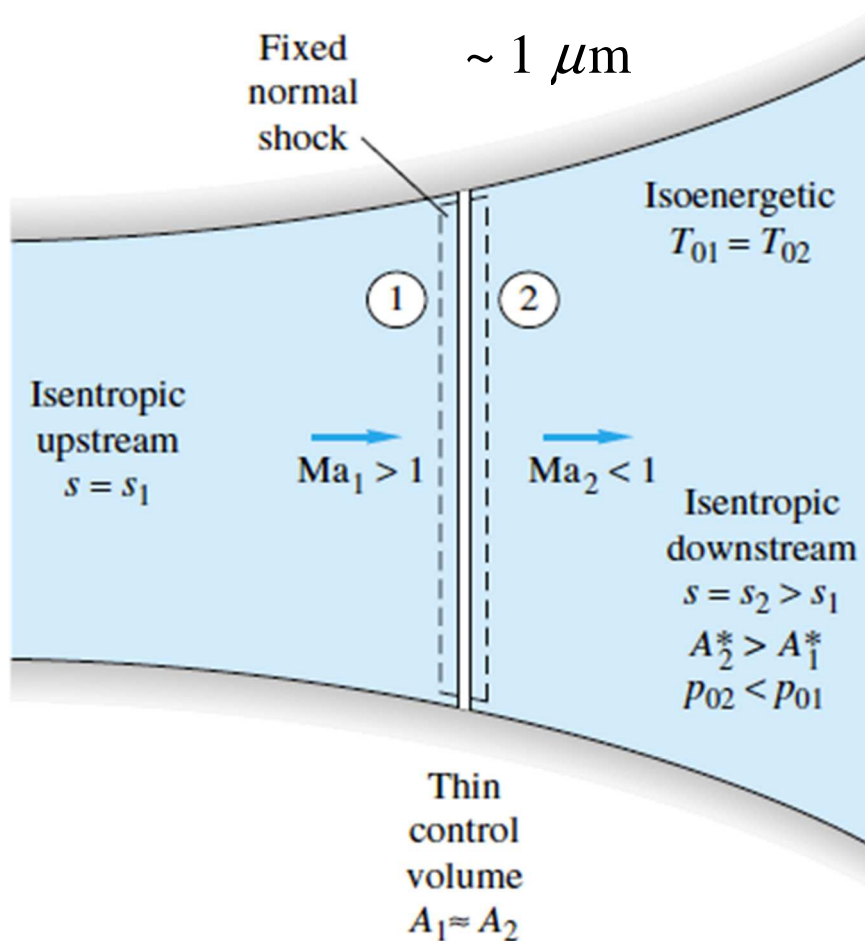
$$\xrightarrow{\kappa = 1.4} \dot{m}_{\max} = \frac{0.6847 p_0 A^*}{(RT_0)^{1/2}}$$

# Example 9.4

Air flows isentropically through a duct. At (1) the area is  $0.05 \text{ m}^2$  and  $V_1 = 180 \text{ m/s}$ ,  $p_1 = 500 \text{ kPa}$ , and  $T_1 = 470 \text{ K}$ . Compute (a)  $T_0$ , (b)  $Ma_1$ , (c)  $p_0$ , and (d) both  $A^*$  and  $\dot{m}$ . If at (2) the area is  $0.036 \text{ m}^2$ , compute  $Ma_2$  and  $p_2$  for (e) subsonic and (f) supersonic flow. Assume  $\kappa = 1.4$ .



# Normal shock wave



continuity:

$$A_1 \approx A_2$$

$$\rho_1 V_1 = \rho_2 V_2 = \text{const.}$$

momentum:

$$p_1 - p_2 = \rho_2 V_2^2 - \rho_1 V_1^2$$

energy:

$$h_1 + \frac{1}{2} \rho_1 V_1^2 = h_2 + \frac{1}{2} \rho_2 V_2^2 = h_0 = \text{const.}$$

ideal gas law:

$$\frac{p_1}{\rho_1 T_1} = \frac{p_2}{\rho_2 T_2} \quad h = c_p T \quad \kappa = \text{const.}$$

# Rankine-Hugueniot relations

eliminate  $V_1$  and  $V_2$  :

$$h_2 - h_1 = \frac{1}{2}(p_2 - p_1) \left( \frac{1}{\rho_2} + \frac{1}{\rho_1} \right)$$

introduce ideal gas law:

$$\frac{\rho_2}{\rho_1} = \frac{1 + \beta p_2/p_1}{\beta + p_2/p_1}, \quad \beta = \frac{\kappa + 1}{\kappa - 1}$$

entropy change accross shock:

$$s_2 - s_1 = c_v \ln \left[ \frac{p_2}{p_1} \left( \frac{\rho_1}{\rho_2} \right)^\kappa \right]$$

2nd Law of Thermodynamics:

$$s_2 \geq s_1 \Rightarrow \begin{cases} p_2 \geq p_1 \\ \rho_2 \geq \rho_1 \end{cases}$$

isentropic flow (Poisson relation):

$$\frac{\rho_2}{\rho_1} = \left( \frac{p_2}{p_1} \right)^{\frac{1}{\kappa}}$$

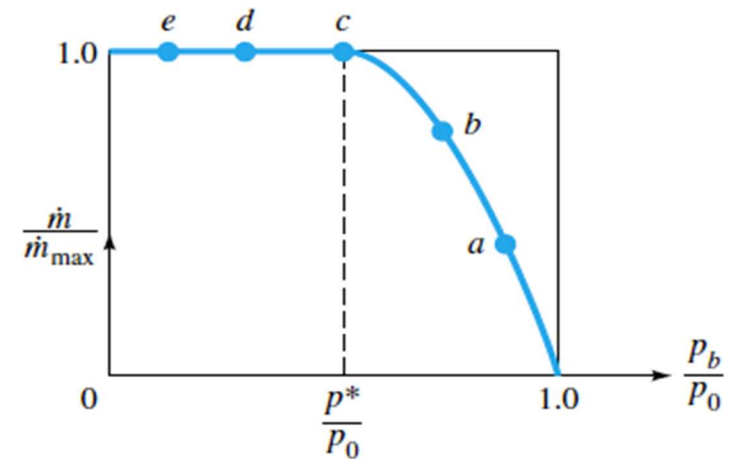
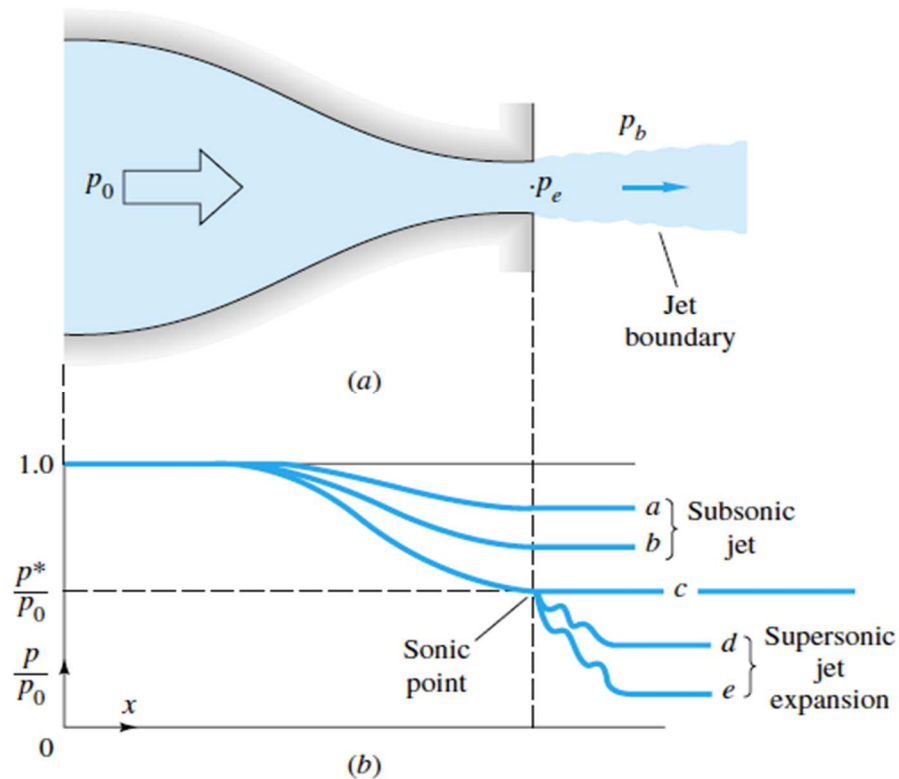
$\frac{p_2}{p_1}$	$\rho_2/\rho_1$		$\frac{s_2 - s_1}{c_v}$
	Eq. (9.51)	Isentropic	
0.5	0.6154	0.6095	-0.0134
0.9	0.9275	0.9275	-0.00005
1.0	1.0	1.0	0.0
1.1	1.00704	1.00705	0.00004
1.5	1.3333	1.3359	0.0027
2.0	1.6250	1.6407	0.0134

# Summary

- $Ma_1 > 1, Ma_2 < 1$
- $\rho_2 > \rho_1$
- $s_2 > s_1, A_2^* > A_1^*$
- weak shocks are nearly isentropic

# Example 9.8

A converging nozzle has a throat area of  $6 \text{ cm}^2$  and stagnation air conditions of  $120 \text{ kPa}$  and  $400 \text{ K}$ . Compute the exit pressure and mass flow if the back pressure is (a)  $90 \text{ kPa}$  and (b)  $45 \text{ kPa}$ . Assume  $\kappa = 1.4$ .



# Source

1. An F/A-18 Hornet photographed just as it broke the sound barrier, photo courtesy of Ensign John Gay, USS Constellation, US Navy,  
The rest of the pictures are from the book of Frank M. White, *Fluid Mechanics*, McGraw-Hill Series in Mechanical Engineering.