

Thermodynamica 2

Thermodynamic relations of systems in equilibrium

Thijs J.H. Vlugt

Engineering Thermodynamics
Process and Energy Department

Lecture 4

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Today:

- meetbare grootheden: C_p , C_v 11.4, 11.5
- Uitzettingscoefficient, compressibiliteit 11.5
- Veranderingen van U , H , A , G met V , P , T 11.6, 11.7
- Homework: Exercises Moran&Shapiro 11.36, 11.37, 11.50 en extra opgaven 2, 5, 7, 8 (zie blackboard)
- Print de laatste 2 slides uit en neem ze mee naar het volgende college
- Volgende week: 11.1, 11.7

Summary

- Partial derivatives

$$df = \left(\frac{\partial f}{\partial x} \right)_y dx + \left(\frac{\partial f}{\partial y} \right)_x dy = Mdx + Ndy \quad \text{so} \quad M = \left(\frac{\partial f}{\partial x} \right)_y \quad \text{en} \quad N = \left(\frac{\partial f}{\partial y} \right)_x$$

- Maxwell Relations

$$\text{If } df = Mdx + Ndy \quad \text{then} \quad \left(\frac{\partial M}{\partial y} \right)_x = \left(\frac{\partial N}{\partial x} \right)_y \quad \text{in other words} \quad \frac{\partial}{\partial x} \left[\left(\frac{\partial f}{\partial y} \right)_x \right]_y = \frac{\partial}{\partial y} \left[\left(\frac{\partial f}{\partial x} \right)_y \right]_x$$

- Integration

$$f(x_2, y_1) = f(x_1, y_1) + \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial x} \right)_y dx$$

- Minus 1 rule (Moran&Shapiro page 495)

$$\left(\frac{\partial x}{\partial y} \right)_z \left(\frac{\partial z}{\partial x} \right)_y \left(\frac{\partial y}{\partial z} \right)_x = -1$$

Sommetje vorige week

- Laat zien dat voor een ideaal gas, U niet afhangt van V , maar dat voor een van der Waals gas U wel afhangt van V

$$\left(\frac{\partial U}{\partial V}\right)_{T,N} = -p + T \left(\frac{\partial p}{\partial T}\right)_{V,N}$$

ideaal gas: $\left(\frac{\partial p}{\partial T}\right)_{V,N} = \frac{nR}{V}$ dus $\left(\frac{\partial U}{\partial V}\right)_{T,N} = -p + \frac{nRT}{V} = 0$

van der Waals gas: $\left(p + \frac{n^2 a}{V^2}\right)(V - nb) = nRT$ oftewel $p = \frac{nRT}{V - nb} - \frac{n^2 a}{V^2}$

invullen: $\left(\frac{\partial p}{\partial T}\right)_{V,N} = \frac{nR}{V - nb}$ dus $\left(\frac{\partial U}{\partial V}\right)_{T,N} = -p + \frac{nRT}{V - nb} = \frac{n^2 a}{V^2}$

NB indien $V \rightarrow \infty$ krijgen we hetzelfde resultaat als voor het ideale gas

Exercise 1

Gegeven is de van der Waals equation of state. Hangt C_v af van het volume?

$$\left(p + \frac{n^2 a}{V^2} \right) (V - nb) = nRT$$

$$dU = TdS - pdV + \mu dN$$

$$dH = TdS + Vdp + \mu dN$$

$$dA = -SdT - pdV + \mu dN$$

$$dG = -SdT + Vdp + \mu dN$$

Exercise 1

Voor elke stof geldig

$$\left(\frac{\partial C_V}{\partial V} \right)_{T,N} = \left(\frac{\partial}{\partial V} \left(\frac{\partial U}{\partial T} \right)_{V,N} \right)_{T,N} = \left(\frac{\partial}{\partial T} \left(\frac{\partial U}{\partial V} \right)_{T,N} \right)_{V,N} = \left(\frac{\partial}{\partial T} \left[-p + T \left(\frac{\partial p}{\partial T} \right)_{V,N} \right] \right)_{V,N} = T \left(\frac{\partial^2 p}{\partial T^2} \right)_{V,N}$$

$$\left(\frac{\partial p}{\partial T} \right)_{V,N} = \frac{R}{v-b} \quad \text{en} \quad \left(\frac{\partial^2 p}{\partial T^2} \right)_{V,N} = 0 \quad \text{dus} \quad \left(\frac{\partial C_V}{\partial V} \right)_{T,N} = 0$$

Invullen voor het systeem dat we bekijken

Partial derivatives of U

$$\left(\frac{\partial U}{\partial T} \right)_{V,N} = C_V$$

$$\left(\frac{\partial U}{\partial V} \right)_{T,N} = -p + T \left(\frac{\partial p}{\partial T} \right)_{V,N}$$

maar wat doen we als we $\left(\frac{\partial U}{\partial p} \right)_{T,N}$ willen weten ?????

Exercise 2

Leid een algemene uitdrukking af voor $C_p - C_v$ als functie van meetbare grootheden.

Bepaal $C_p - C_v$ voor een gas met de volgende toestandsvergelijking:

$$v = RT/p - A/T + B$$

Exercise 2 (....continued)

$$C_p - C_V = \left(\frac{\partial H}{\partial T} \right)_{p,N} - \left(\frac{\partial U}{\partial T} \right)_{V,N} = \left(\frac{\partial U}{\partial T} \right)_{p,N} + p \left(\frac{\partial V}{\partial T} \right)_{p,N} - \left(\frac{\partial U}{\partial T} \right)_{V,N}$$

$H = U + pV$

$$dU = \left(\frac{\partial U}{\partial T} \right)_{V,N} dT + \left(\frac{\partial U}{\partial V} \right)_{T,N} dV$$

$U(T, V)$

$$dV = \left(\frac{\partial V}{\partial T} \right)_{p,N} dT + \left(\frac{\partial V}{\partial p} \right)_{T,N} dp$$

$V(T, p)$

$$dU = \left(\frac{\partial U}{\partial T} \right)_{V,N} dT + \left(\frac{\partial U}{\partial V} \right)_{T,N} \left(\frac{\partial V}{\partial T} \right)_{p,N} dT + \left(\frac{\partial U}{\partial V} \right)_{T,N} \left(\frac{\partial V}{\partial p} \right)_{T,N} dp$$

invullen

$$\left(\frac{\partial U}{\partial T} \right)_{p,N} = \left(\frac{\partial U}{\partial T} \right)_{V,N} + \left(\frac{\partial U}{\partial V} \right)_{T,N} \left(\frac{\partial V}{\partial T} \right)_{p,N}$$

p constant ($dp=0$) en
delen door dT

$$\left(\frac{\partial U}{\partial V} \right)_{T,N} = -p + T \left(\frac{\partial p}{\partial T} \right)_{V,N}$$

eerder afgeleid

$$C_p - C_V = T \left(\frac{\partial p}{\partial T} \right)_{V,N} \left(\frac{\partial V}{\partial T} \right)_{p,N}$$

uiteindelijke resultaat, geldig
voor ieder systeem

Exercise 2

$$C_p - C_V = T \left(\frac{\partial p}{\partial T} \right)_{V,N} \left(\frac{\partial V}{\partial T} \right)_{p,N}$$

Geldig voor elk systeem

$$\left(\frac{\partial p}{\partial T} \right)_{V,N} = \frac{R(V - B) + \frac{2RA}{T}}{\left(V + \frac{A}{T} - B \right)^2}$$

Invullen voor het gevraagde systeem

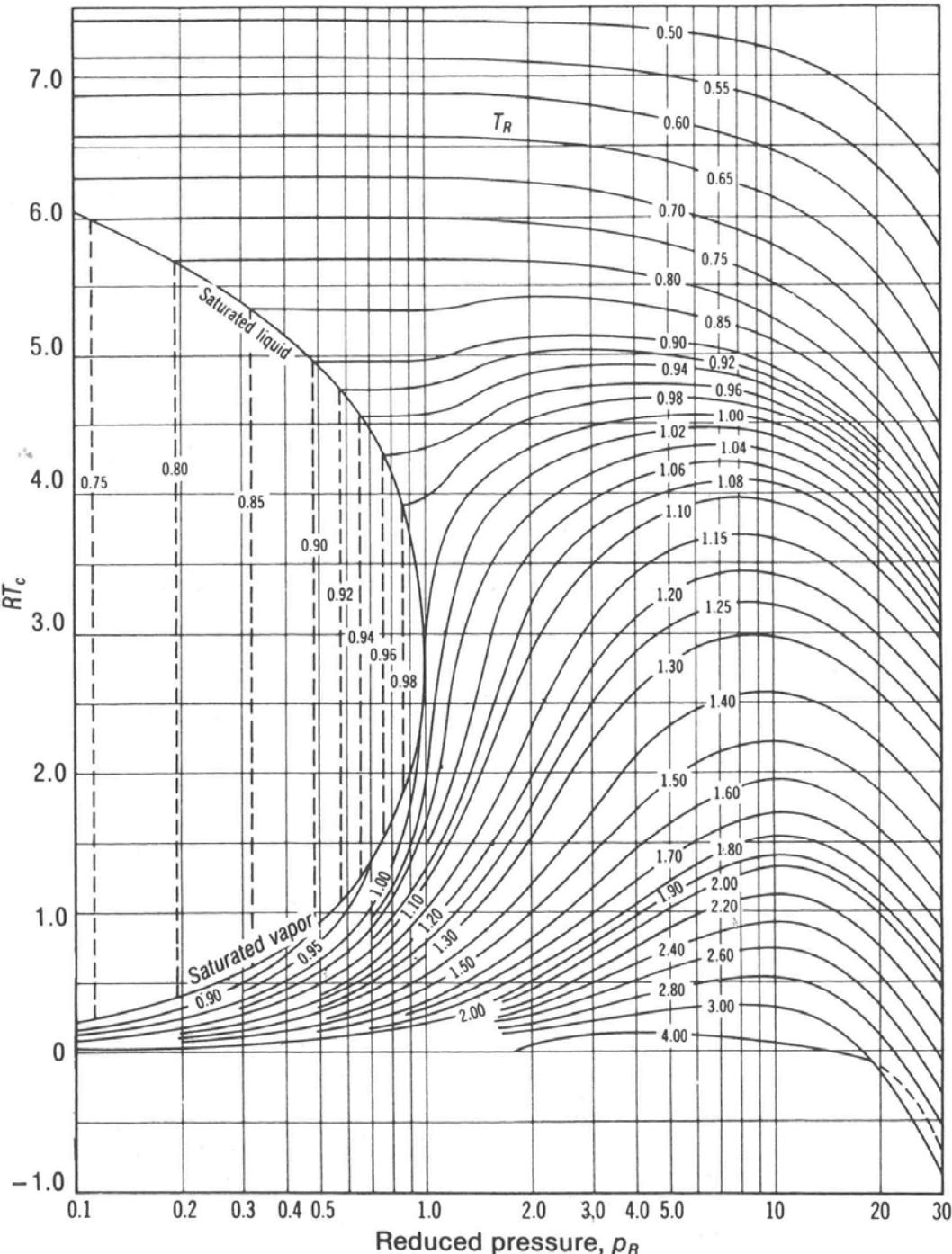
$$\left(\frac{\partial V}{\partial T} \right)_{p,N} = \frac{R}{p} + \frac{A}{T^2}$$
$$C_p - C_V = T \left(\frac{\partial p}{\partial T} \right)_{V,N} \left(\frac{\partial V}{\partial T} \right)_{p,N} = T \left[\frac{R}{p} + \frac{A}{T^2} \right] \left[\frac{R(V - B) + \frac{2RA}{T}}{\left(V + \frac{A}{T} - B \right)^2} \right]$$

Enthalpy departure in corresponding states

caution

$\left(\frac{\bar{h} - \bar{h}^{ig}}{\bar{R}T_c} \right)$ is equal to minus the values from this table

asterisk * (just as *ig*) is for ideal gas



Entropy-departure in corresponding states

