

Microscopic theory of superconductivity

- Superfluidity
- Phonon attraction
- Cooper pairs
- BCS theory
- Energy gap
- Correlation length
- Type-I and type-II superconductors

Superfluidity

Liquid Helium-4: at low temperature flows with zero viscosity

Relative to the liquid (velocity \mathbf{v}): viscosity means creation of quasiparticles

Momentum \mathbf{p} ; energy $\varepsilon(\vec{p})$

Relative to the ground: momentum $\vec{P}' = \vec{P} + M\vec{v}$
energy $E' = E + \vec{P}\vec{v} + Mv^2/2$
 M – mass of the liquid

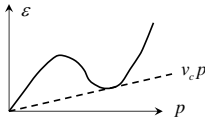
Energy change associated with the creation of a quasiparticle:

$$\Delta E = \varepsilon(\vec{p}) + \vec{p}\vec{v} \quad \text{If it is positive, superfluid (no viscosity)}$$

$$v < v_c = \min(\varepsilon(p)/p)$$

Superfluidity

Helium-4
(bosons)



Fermi-system: Elementary excitations are electron-hole pairs
Energy: may be arbitrarily small
Momentum: can be finite

No superfluidity!

Maybe electrons are bound in pairs and behave like Helium atoms.
Difficulty: Why do not they repel each other then?

Electron-phonon interaction again

One electron emits a phonon, another one absorbs this phonon:
Mechanism for electron-electron interaction!

Initial state: electron #1 with \vec{p}_1 and electron #2 with \vec{p}_2

Intermediate state: electron #1 with $\vec{p}_1 - \hbar\vec{k}$, electron #2 with \vec{p}_2
phonon with \vec{k}

Final state: electron #1 with $\vec{p}_1 - \hbar\vec{k}$ and electron #2 with $\vec{p}_2 + \hbar\vec{k}$

$$\varepsilon(\vec{p}_1) + \varepsilon(\vec{p}_2) = \varepsilon(\vec{p}_1 - \hbar\vec{k}) + \varepsilon(\vec{p}_2 + \hbar\vec{k}) \quad \text{-same energy of initial and final states}$$

Second-order perturbation theory: $\Delta E_i = \sum_r \frac{|M_{ir}|^2}{E_i - E_r}$

$$\Delta E_{1\vec{p}_1} = \frac{|V_{\vec{k}}|^2}{\varepsilon(\vec{p}_1) - \varepsilon(\vec{p}_1 - \hbar\vec{k}) - \hbar\omega(\vec{k})} = \frac{\lambda}{gV} \frac{\hbar\omega(\vec{k})}{\varepsilon(\vec{p}_1) - \varepsilon(\vec{p}_1 - \hbar\vec{k}) - \hbar\omega(\vec{k})}$$

Electron-phonon interaction again

Another process contributing: electron absorbs a phonon with $-\mathbf{k}$

Initial state: electron #1 with \vec{p}_1 and electron #2 with \vec{p}_2

Intermediate state: electron #1 with \vec{p}_1 , electron #2 with $\vec{p}_2 + \hbar\vec{k}$
phonon with $-\vec{k}$

Final state: electron #1 with $\vec{p}_1 - \hbar\vec{k}$ and electron #2 with $\vec{p}_2 + \hbar\vec{k}$

$$\varepsilon(\vec{p}_1) + \varepsilon(\vec{p}_2) = \varepsilon(\vec{p}_1 - \hbar\vec{k}) + \varepsilon(\vec{p}_2 + \hbar\vec{k}) \quad \text{-same energy of initial and final states}$$

Second-order perturbation theory: $\Delta E_i = \sum_r \frac{|M_{ir}|^2}{E_i - E_r}$

$$\Delta E_{2\vec{p}_1} = \frac{|V_{\vec{k}}|^2}{\varepsilon(\vec{p}_2) - \varepsilon(\vec{p}_2 + \hbar\vec{k}) - \hbar\omega(\vec{k})} = -\frac{\lambda}{gV} \frac{\hbar\omega(\vec{k})}{\varepsilon(\vec{p}_1) - \varepsilon(\vec{p}_1 - \hbar\vec{k}) + \hbar\omega(\vec{k})}$$

Phonon attraction

Summing up:

$$\Delta E_{\vec{p}_1} = \Delta E_{1\vec{p}_1} + \Delta E_{2\vec{p}_1} = \frac{2\lambda}{gV} \frac{[\hbar\omega(\vec{k})]^2}{[\varepsilon(\vec{p}_1) - \varepsilon(\vec{p}_1 - \hbar\vec{k})]^2 - [\hbar\omega(\vec{k})]^2}$$

$$\text{For } \varepsilon(\vec{p}_1) - \varepsilon(\vec{p}_1 - \hbar\vec{k}) \ll \hbar\omega(\vec{k}) \quad \Delta E = -\frac{2\lambda}{gV}$$

Corresponds to attracting point-like interaction!!

$$V_{ph} = -\frac{2\lambda}{g} \delta(\vec{r} - \vec{r}')$$

Phonon attraction

Electron-electron attraction:
$$V_{ph} = -\frac{2\lambda}{g} \delta(\vec{r} - \vec{r}')$$

Compare this with Coulomb repulsion:

$$V_C = -\frac{e^2}{|\vec{r} - \vec{r}'|} \Rightarrow e^2 a_B^2 \delta(\vec{r} - \vec{r}')$$

$$\begin{aligned} g &\sim mp_F / \hbar^3 \\ a_B &\sim \hbar^2 / me^2 \end{aligned}$$

$$\frac{V_{ph}}{V_C} \sim \frac{\lambda}{ge^2 a_B^2} \sim \frac{e^2 \lambda}{\hbar v_F} \sim 1 \quad \text{Same order of magnitude}$$

Attraction dominates > superconductivity
Repulsion dominates > no superconductivity

Cooper pairs

We found a (weak) electron-electron attraction, but it does not mean two electrons form a bound state (a pair).

Quantum mechanics: an electron in a shallow quantum well:

1D: always a bound state

2D and 3D: no bound state (only if the well is deep enough)

Now, we have not free electrons, but quasiparticles above the Fermi surface.

Two-particle Schrödinger equation:

$$[\hat{H}_1 + \hat{H}_2 + U(\vec{r}_1, \vec{r}_2)] \Psi(\vec{r}_1, \vec{r}_2) = E \Psi(\vec{r}_1, \vec{r}_2)$$

Ground state: total momentum and total spin of the pair equal zero

$$\Psi(\vec{r}_1, \vec{r}_2) = \sum_{\vec{p}} c_{\vec{p}} \psi_{\vec{p}\uparrow}(\vec{r}_1) \psi_{-\vec{p}\downarrow}(\vec{r}_2)$$

Cooper pairs

Back to Schrödinger equation:

$$2\xi_{\vec{p}} c_{\vec{p}} + \sum_{\vec{p}'} U_{\vec{p}\vec{p}'} c_{\vec{p}'} = E c_{\vec{p}}$$

$$U_{\vec{p}\vec{p}'} = -U_0, p_F - \frac{\hbar\omega_D}{v_F} < p, p' < p_F + \frac{\hbar\omega_D}{v_F}$$

Solution: $E = -2\Delta$

$$c_{\vec{p}} = \frac{U_0}{2|\xi_{\vec{p}} + \Delta|} \sum_{\vec{p}'} c_{\vec{p}'} \Rightarrow \frac{U_0 g(\epsilon_F)}{2} \ln\left(\frac{\hbar\omega_D}{\Delta}\right) = 1$$

$$\Delta = \hbar\omega_D \exp(-2/g(\epsilon_F)U_0) = \hbar\omega_D \exp(-4/\lambda)$$

No small parameter – non-perturbative calculation.

BCS Hamiltonian

$$\hat{H} = \sum_{\vec{p}\sigma} \xi_{\vec{p}} \hat{c}_{\vec{p}\sigma}^\dagger \hat{c}_{\vec{p}\sigma} - U_0 \sum_{\vec{p}_1 + \vec{p}_2 = \vec{p}_1 + \vec{p}_2} \hat{c}_{\vec{p}_1\uparrow}^\dagger \hat{c}_{\vec{p}_2\downarrow}^\dagger \hat{c}_{\vec{p}_2\downarrow} \hat{c}_{\vec{p}_1\uparrow}$$

Kinetic energy

Phonon attraction (only electrons with different spin interact)

What is the ground state corresponding to this Hamiltonian?
(Advanced quantum mechanics, Lecture 4)

$$\Psi = \prod_{\vec{p}} [u_{\vec{p}} + v_{\vec{p}} \hat{c}_{\vec{p}\uparrow}^\dagger \hat{c}_{-\vec{p}\downarrow}] \quad |u_{\vec{p}}|^2 + |v_{\vec{p}}|^2 = 1 \quad \text{Normal state: } v=0$$

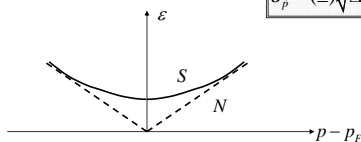
Need to find u and v from the minimisation of energy (variational principle)

BCS theory

$$u_{\vec{p}} = \frac{1}{\sqrt{2}} \sqrt{1 + \xi_{\vec{p}} / \epsilon_{\vec{p}}}, v_{\vec{p}} = \frac{1}{\sqrt{2}} \sqrt{1 - \xi_{\vec{p}} / \epsilon_{\vec{p}}}$$

Energy spectrum of quasiparticles:

$$\epsilon_{\vec{p}} = (\pm) \sqrt{\Delta^2 + \xi_{\vec{p}}^2}$$

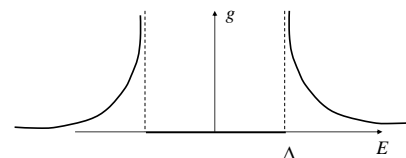


No quasiparticles with energy below Δ !

Density of states

$$\epsilon_{\vec{p}} = \sqrt{\Delta^2 + \xi_{\vec{p}}^2}$$

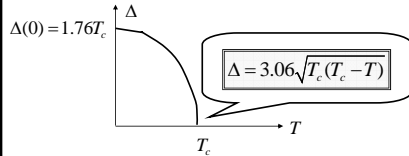
$$g(E) = \sum_{\vec{p}} \delta(E - \sqrt{\Delta^2 + \xi_{\vec{p}}^2}) \propto \frac{\theta(E - \Delta)}{\sqrt{E - \Delta}}$$



Energy gap

$$T_c = 1.4 \hbar \omega_p U \exp(-2/(gU_0)) \quad \text{- transition temperature}$$

Energy spectrum of quasiparticles:



Energy gap can be chosen as the order parameter for the Landau description of the second order phase transition – see next Lecture.

Correlation length

What is a typical size of a Cooper pair?

$$\delta p \sim \Delta / v \implies \delta r \sim \hbar v / \Delta$$

Example: $\Delta / k_B \sim T_c \sim 10K \implies \delta r \sim 10^{-6} m$
 $v \sim 10^6 m/s$

The "size" of a Cooper pair is much bigger than the distance between electrons: It does not make any sense.

It is better to talk about the correlation between electrons.

Correlation length: $\xi(T) \sim \hbar v / \Delta(T)$ Diverges at the transition point!

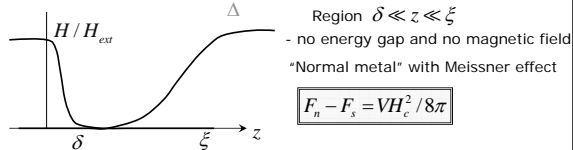
London equations: correspond to $\xi=0$.

Pippard limit

Correlation length: $\xi(T) \sim \hbar v / \Delta(T)$ Both diverge as $1/\sqrt{T_c - T}$

Penetration depth: $\delta = \sqrt{mc^2 / (4\pi n_s e^2)}$

Assume $\delta \ll \xi$ and consider a boundary between N and S



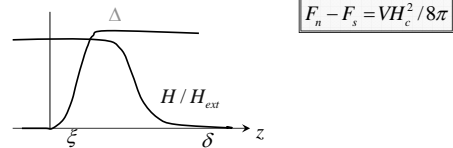
Positive boundary energy -> intermediate state

London limit

Assume now $\delta \gg \xi$

Region $\xi \ll z \ll \delta$ - both energy gap and magnetic field present.

"Superconductor" without Meissner effect



Negative boundary energy -> unusual behavior of a superconductor in an external field

Two types of superconductors

Type I: positive boundary energy $\delta \ll \xi$

Type II: negative boundary energy $\delta \gg \xi$

Criterion: $\kappa(T) = \delta(T) / \xi(T)$ -only weakly depends on temperature

$$\kappa < 1/\sqrt{2} \quad \text{- type I}$$

$$\kappa > 1/\sqrt{2} \quad \text{- type II}$$

Bogoliubov – de Gennes equations

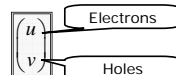
Mean-field treatment of BCS Hamiltonian

$$\hat{H} = \sum_{\vec{p}\sigma} \xi_{\vec{p}} \hat{c}_{\vec{p}\sigma}^\dagger \hat{c}_{\vec{p}\sigma} + \Delta \sum_{\vec{p}} \hat{c}_{\vec{p}\uparrow}^\dagger \hat{c}_{-\vec{p}\downarrow}^\dagger + \Delta^* \sum_{\vec{p}} \hat{c}_{\vec{p}\downarrow} \hat{c}_{-\vec{p}\uparrow}$$

It is quadratic: can be diagonalized

$$\psi_{\vec{p}\uparrow} = u_{\vec{p}} \hat{c}_{\vec{p}\uparrow} + v_{\vec{p}} \hat{c}_{-\vec{p}\downarrow}^\dagger \quad \text{Not the same } u \text{ and } v \text{ as in BCS!}$$

Two-component wave function:



Bogoliubov – de Gennes equations

Minimizing the Hamiltonian:
Bogoliubov – de Gennes equations – Two-component generalization
of Schrodinger equation

Best in the coordinate representation:

$$\begin{cases} \hat{H}u + \Delta v = \varepsilon u \\ -\hat{H}v + \Delta^* u = \varepsilon v \end{cases}$$

Energy spectrum: same as BCS

$$\varepsilon_{\vec{p}} = \pm \sqrt{\Delta^2 + \xi_{\vec{p}}^2}$$