Offshore Hydromechanics

Module 1 :

Hydrostatics Constant Flows Surface Waves

OE4620 Offshore Hydromechanics

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Offshore Engineering



Delft University of Technology



First hour:

- Schedule for remainder of hydromechanics module 1
- Assignment stability
- Basic flow properties
- Potential flow concepts

Second hour:

- Potential flow elements
- Superposition of flow elements
- Cylinder in uniform flow



Schedule OE4620 Module 1

- 13 November Constant potential flow phenomena
- 20 November Constant real flow phenomena 1
- 27 November Constant real flow phenomena 2
- 4 December Ocean surface waves 1 (!)
- 11 December Ocean surface waves 2
- 18 December Extra/Examples



Assignment: Stability

- Make the assignment for next lecture
- Next lecture the assignment will be worked out briefly
- Questions regarding the assignment or the subjects of module 1? Contact:
 - Room 2.78
 - Tel. 015 278 7568
 - E-mail w.e.devries@tudelft.nl



CONSTANT POTENTIAL FLOW PHENOMENA

Flows in this chapter:

- Obey simplified laws of fluid mechanics
- Are subject to limitations
- Give qualitative impressions of flow phenomena



Basic Flow Properties

- Ideal fluid :
 - Non viscous
 - Incompressible
 - Continuous
 - Homogeneous



•Increase of mass per unit time:

$$\frac{\partial m}{\partial t} = -\frac{\partial}{\partial t} \left(\rho \cdot dx dy dz \right)$$





•Mass through a plane *dxdz* during a unit of time:

 $m_{in} = \rho \cdot v \cdot dx dz dt$





•Mass through a plane *dxdz* during a unit of time:

$$m_{in} = \rho \cdot v \cdot dx dz dt$$
$$m_{out} = \left[\rho v + \frac{\partial (\rho v)}{\partial y} dy\right] \cdot dx dz dt$$













Combining for *x*, *y* and *z* direction yields the **Continuity Equation**

 $\frac{\partial m}{\partial t}$



Combining for *x*, *y* and *z* direction yields the **Continuity Equation**

$$\frac{\partial m}{\partial t} = -\frac{\partial \rho}{\partial t} dx dy dz$$



Combining for *x*, *y* and *z* direction yields the **Continuity Equation**

$$\frac{\partial m}{\partial t} = -\frac{\partial \rho}{\partial t} dx dy dz = \left\{ \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} \right\} dx dy dz$$



Combining for *x*, *y* and *z* direction yields the **Continuity Equation**

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} = 0$$



For an ideal, *incompressible* fluid:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$



For an ideal, *incompressible* fluid:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

Or, with
$$\nabla = \frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z}$$
 : $\nabla \overline{V} = 0$



Deformation

$$\frac{\partial v}{\partial x} = \tan \dot{\alpha} \approx \dot{\alpha} \qquad \frac{\partial u}{\partial y} = \tan \dot{\beta} \approx \dot{\beta}$$





Deformation

$$\frac{\partial v}{\partial x} = \tan \dot{\alpha} \approx \dot{\alpha} \qquad \frac{\partial u}{\partial y} = \tan \dot{\beta} \approx \dot{\beta}$$

Deformation velocity (dilatation):

$$\frac{\dot{\alpha} + \dot{\beta}}{2} = \frac{1}{2} \left(\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$





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Rotation (in 2 D) :

$$\dot{\varphi} = \frac{1}{2} \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$





Velocity Potential

A scalar function, Φ

The velocity in any point of the fluid and in any direction is the derivative to that direction of the potential

$$u = \frac{\partial \Phi}{\partial x}$$
 $v = \frac{\partial \Phi}{\partial y}$ $w = \frac{\partial \Phi}{\partial z}$



Properties of Velocity Potentials

• Combining the continuity condition with the velocity potentials results in the *Laplace* equation :

$$\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0 \quad \text{or} \quad \nabla^2 \Phi = 0$$

• Potential flow is rotation-free by definition :

$$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0 \qquad \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} = 0 \qquad \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} = 0$$



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• Euler: Apply Newton's 2nd law to non viscous and incompressible fluid

$$dm \frac{Du}{Dt}$$



• Euler: Apply Newton's 2nd law to non viscous and incompressible fluid

$$dm\frac{Du}{Dt} = \rho \cdot dxdydz \cdot \frac{Du}{Dt} =$$



• Euler: Apply Newton's 2nd law to non viscous and incompressible fluid

$$dm\frac{Du}{Dt} = \rho \cdot dxdydz \cdot \frac{Du}{Dt} = -\frac{\partial p}{\partial x}dx \cdot dydz$$

Mass* acceleration = pressure*area



• Euler Equations:

$$\frac{\partial u}{\partial t} = u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial x}$$

 $\left(\frac{\partial \Phi}{\partial y}\right)^2$

With

$$u\frac{\partial u}{\partial x} = \frac{\partial \Phi}{\partial x} \cdot \frac{\partial^2 \Phi}{\partial x^2} = \frac{1}{2} \cdot \frac{\partial}{\partial x} \left(\frac{\partial \Phi}{\partial x}\right)^2$$

$$v\frac{\partial u}{\partial y} = \frac{\partial \Phi}{\partial y} \cdot \frac{\partial^2 \Phi}{\partial x \partial y} = \frac{1}{2} \cdot \frac{\partial}{\partial x} \left(\frac{\partial \Phi}{\partial y}\right)^2$$

$$w\frac{\partial u}{\partial z} = \frac{\partial \Phi}{\partial z} \cdot \frac{\partial^2 \Phi}{\partial x \partial z} = \frac{1}{2} \cdot \frac{\partial}{\partial x} \left(\frac{\partial \Phi}{\partial z}\right)^2$$



• Differentiate w.r.t. *x*, *y* or *z* gives 0:

$$\frac{\partial}{\partial x} \left\{ \frac{\partial \Phi}{\partial t} + \frac{1}{2} \left[\left(\frac{\partial \Phi}{\partial x} \right)^2 + \left(\frac{\partial \Phi}{\partial y} \right)^2 + \left(\frac{\partial \Phi}{\partial z} \right)^2 \right] + \frac{p}{\rho} \right\} = 0 \qquad \text{(similar for } \gamma,$$

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Z)

• Differentiate w.r.t. *x*, *y* or *z* gives 0:

$$\frac{\partial}{\partial x} \left\{ \frac{\partial \Phi}{\partial t} + \frac{1}{2} \left[\left(\frac{\partial \Phi}{\partial x} \right)^2 + \left(\frac{\partial \Phi}{\partial y} \right)^2 + \left(\frac{\partial \Phi}{\partial z} \right)^2 \right] + \frac{p}{\rho} \right\} = 0 \qquad \text{(similar for } y, z\text{)}$$

• This expression is a function of time only, providing the **Bernoulli Equation:**

$$\frac{\partial \Phi}{\partial t} + \frac{1}{2}V^2 + \frac{p}{\rho} + gz = C(t)$$



• For a *stationary* flow along a streamline :

$$\frac{1}{2}V^2 + \frac{p}{\rho} + gz = C$$

• This can be written in two alternative ways :

$$\frac{V^{2}}{2g} + \frac{p}{\rho g} + z = C \qquad \qquad \frac{1}{2}\rho V^{2} + p + \rho gz = C$$



Stream Function Ψ

- Ψ = constant along a streamline
- Definition : $u = \frac{\partial \Psi}{\partial y}$ and $v = -\frac{\partial \Psi}{\partial x}$
- Stream lines and equipotential lines are orthogonal
- No cross flow, impervious boundary of stream tube





Potential Flow Elements

- Uniform flow
- Source
- Sinks
- Circulation

Superposition is allowed



Uniform Flow



$$\Phi = +U \cdot x \qquad \Phi = -U \cdot x$$
$$\Psi = +U \cdot y \qquad \Psi = -U \cdot y$$



Source & Sink



c. Source

d. Sink



 $\Phi = +\frac{Q}{2\pi} \cdot \ln r$ $\Psi = -\frac{Q}{2\pi} \cdot \theta$



Circulation (Vortex)

• Counter clockwise flow:

$$\Phi = + \frac{\Gamma}{2\pi} \cdot \theta \qquad \Psi = -\frac{\Gamma}{2\pi} \cdot \ln r$$

Clockwise flow

$$\Phi = -\frac{\Gamma}{2\pi} \cdot \theta \qquad \Psi = +\frac{\Gamma}{2\pi} \cdot \ln r$$



e. Circulation



















Uniform flow + sink





Separated Source + sink



 $\overline{\mathbf{X}} \quad \Psi = \frac{Q}{2\pi} \cdot \arctan\left(\frac{2ys}{x^2 + y^2 + s^2}\right)$



Uniform flow + source + sink

- No flow through boundary (= streamline)
- Rankine ship forms



$$\Psi = \frac{Q}{2\pi} \cdot \arctan\left(\frac{2ys}{x^2 + y^2 - s^2}\right) + U_{\infty} \cdot y$$



Doublet or Dipole

- A source and sink pair placed very close together
- (s -> 0)





Dipole + uniform flow

- Model for cylinder
- X-axis is a streamline
- Boundary at $R = \sqrt{\frac{\mu}{U_{\infty}}}$

is also a streamline







Uniform flow around a cylinder



$$v_{\theta} = -\left[\frac{\partial\Psi}{\partial r}\right]_{r=R} = -\frac{\partial}{\partial r}\left\{\frac{\mu\sin\theta}{r} - U_{\infty}r\sin\theta\right\}_{r=R} \qquad v_{\theta} = -2U_{\infty}\sin\theta$$

- $V_{\theta} = 0$ at stagnation points
- $V_{\theta} = 2^* U_{\infty}$ at sides of cylinder



Uniform flow + circulation

Pressure asymmetry creates a lift force





Continuity Equation, general case

$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u)}{\partial x} + \frac{\partial (\rho v)}{\partial y} + \frac{\partial (\rho w)}{\partial z} = 0$$

Ideal fluid, incompressible

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$



Deformation velocity or dilatation:

$$\frac{\dot{\alpha} + \dot{\beta}}{2} = \frac{1}{2} \left(\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$

Rotation (in 2 D) :

$$\dot{\varphi} = \frac{1}{2} \left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$



Velocity Potential

A scalar function, Φ

The velocity in any point of the fluid and in any direction is the derivative to that direction of the potential

 $u = \partial \Phi / \partial x$ etc



Properties of Velocity Potentials

• The continuity condition results in the *Laplace* equation :

$$\left|\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2} = 0\right| \qquad \text{or:} \qquad \boxed{\left|\nabla^2 \Phi = 0\right|}$$

• Potential flow is rotation-free by definition :

$$\left|\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0\right|$$



Force = mass * acceleration : Euler non viscous and incompressible fluid

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho} \cdot \frac{\partial p}{\partial x} \bigg|$$

Energy conservation along a flowline : Bernoulli non viscous and incompressible fluid

$$\frac{\partial}{\partial x} \left\{ \frac{\partial \Phi}{\partial t} + \frac{1}{2} \left[\left(\frac{\partial \Phi}{\partial x} \right)^2 + \left(\frac{\partial \Phi}{\partial y} \right)^2 + \left(\frac{\partial \Phi}{\partial z} \right)^2 \right] + \frac{p}{\rho} \right\} = 0 \qquad \qquad \boxed{\frac{\partial \Phi}{\partial t} + \frac{1}{2} V^2 + \frac{p}{\rho} + gz = C(t)}$$

stationary flow :
$$\frac{\left|\frac{1}{2}\rho V^2 + p + \rho gz = Constant\right|}{\left|\frac{1}{2}\rho V^2 + p + \rho gz = Constant\right|}$$



Stream Function Ψ

- Companion of the Velocity Potential
- Definition : $u = \frac{\partial \Psi}{\partial y}$ (a) and $v = -\frac{\partial \Psi}{\partial x}$ (b)
- Stream lines and equipotential lines are orthogonal
- No cross flow, impervious boundary of stream tube





Potential Flow Elements

- Uniform flow
- Sources and sinks
- Circulating flows

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Superposition is allowed
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Polar Coordinate Description



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Uniform flow + sink





Source + sink





Uniform flow + source + sink





Dipole

Strength $\mu = Qs/\pi$





Dipole + uniform flow





Uniform flow around a cylinder



Stagnation points on the X axis $V_{max} = 2 U$ at Y axis



Uniform flow + circulation

Pressure asymmetry creates a lift force



